

# On the Status of the Longitudinal Emittance Measurement at the UNILAC

Timo Milosic

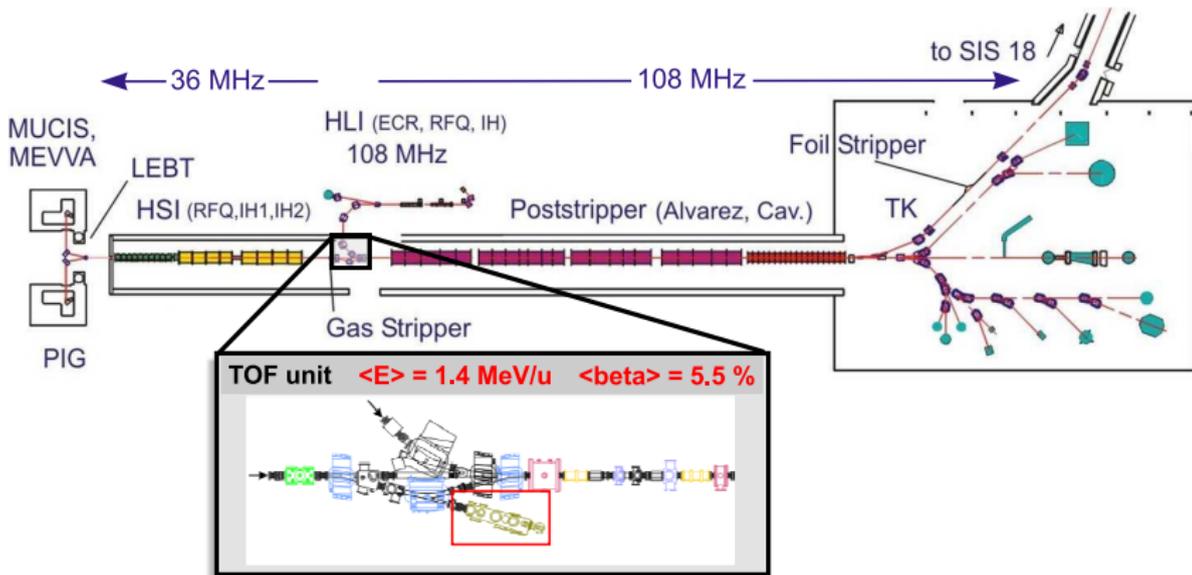
B/SD Group Seminar 02/11/10

February 15, 2010

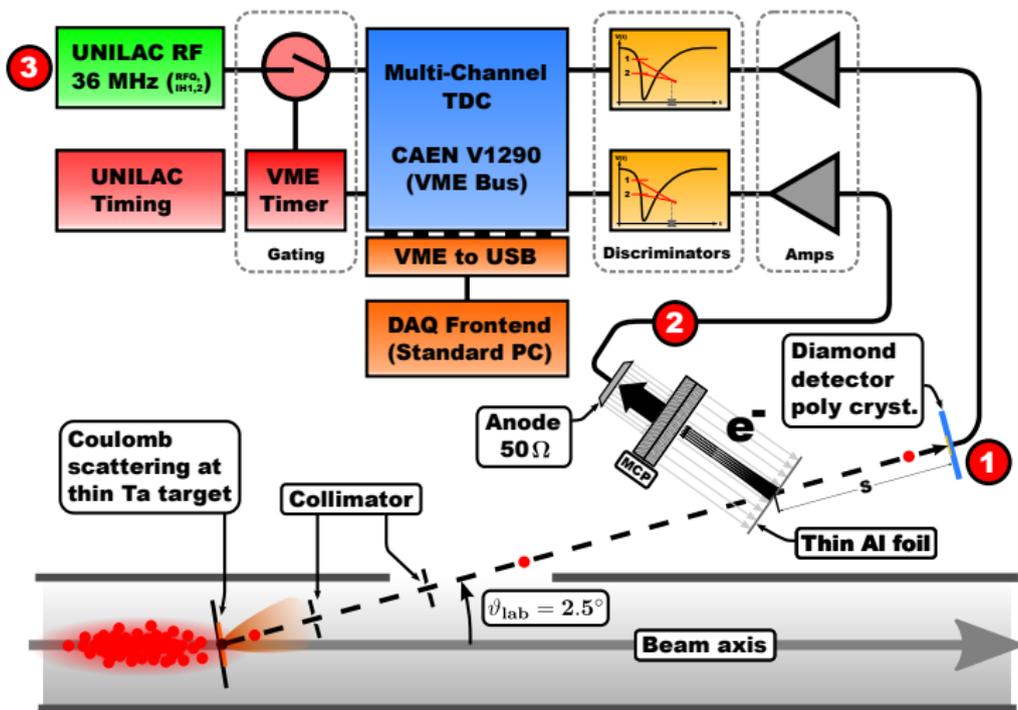
# Outline

- Overview of the measurement setup
- Estimation of systematic errors
- Affect of a limited resolution on measured beam parameters
- Recent measurements
- SC Diamond and Post-Processing

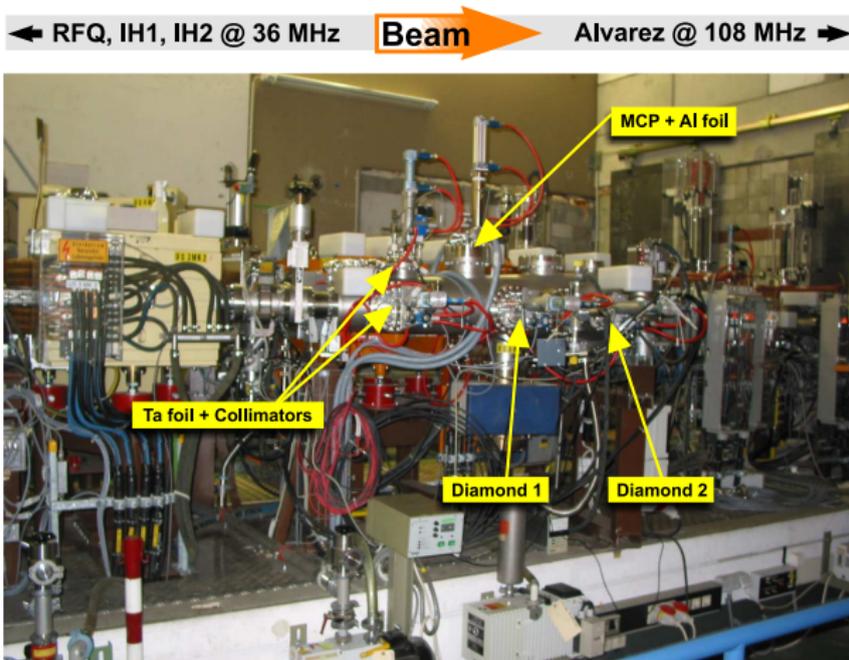
# UNILAC Site



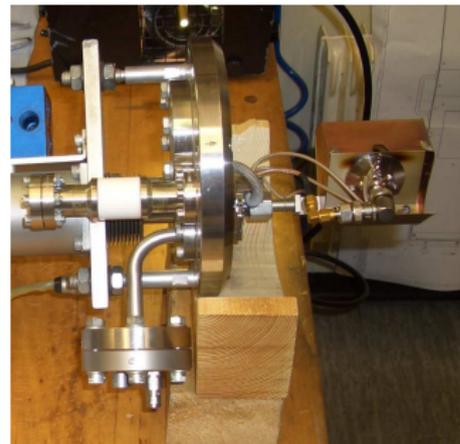
# Schematic operation



# Setup/Installation

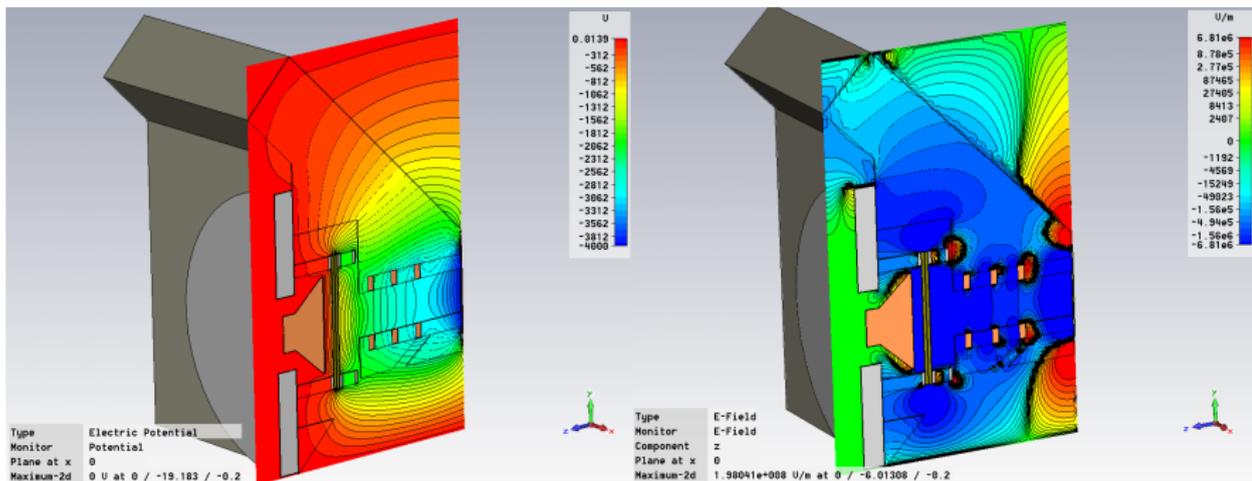


# MCP (Multi Channel Plate)

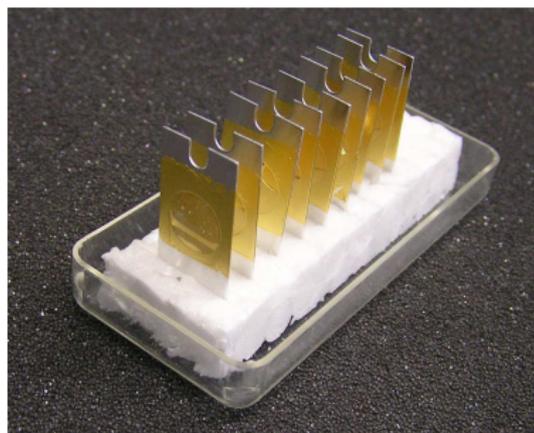
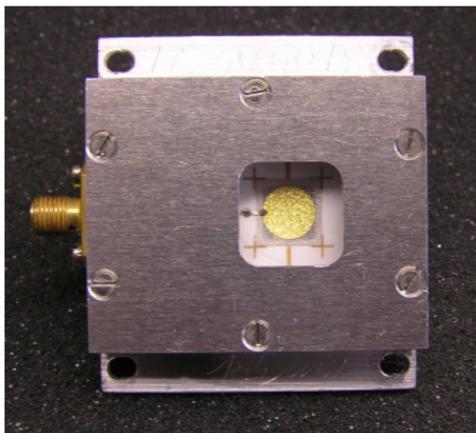


MCP (50  $\Omega$  setup) & assembled MCP actuator

# MCP CST EM Field Simulation

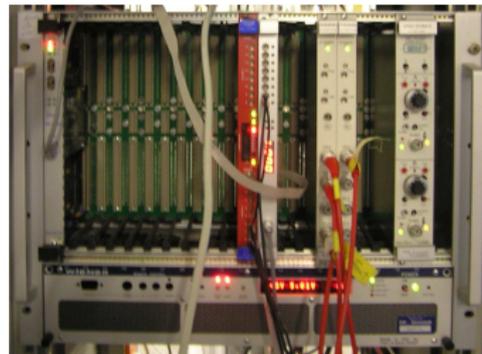
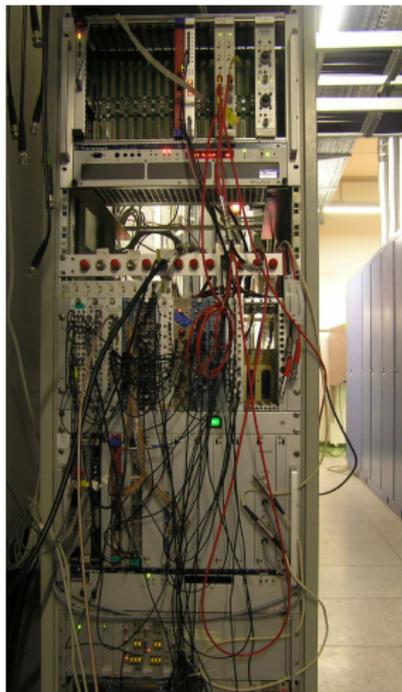


# Foils & Diamond detector



Diamond detector & Au foils (exemplar)

# DAQ



## Redesigned DAQ (2007)

- VME to USB interface, TDC (CAEN V1290), timing module, 3 HV supplies.

# TDC Benchmark

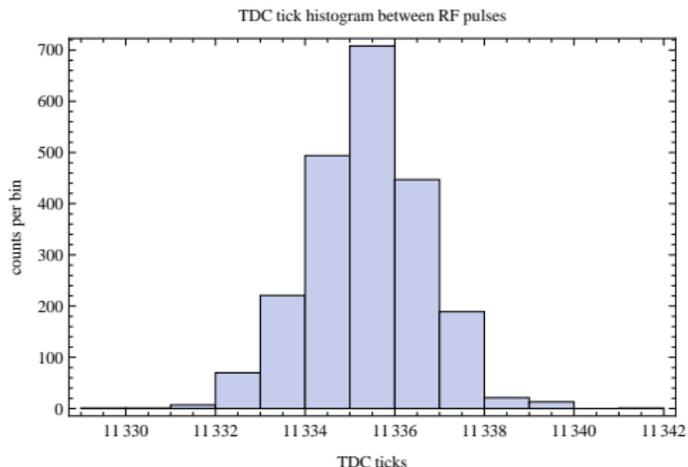
TDC (CAEN V1290):

- LSB of  $\approx 25$  ps.
- RMS resolution 35 ps typical.

Timing tests:

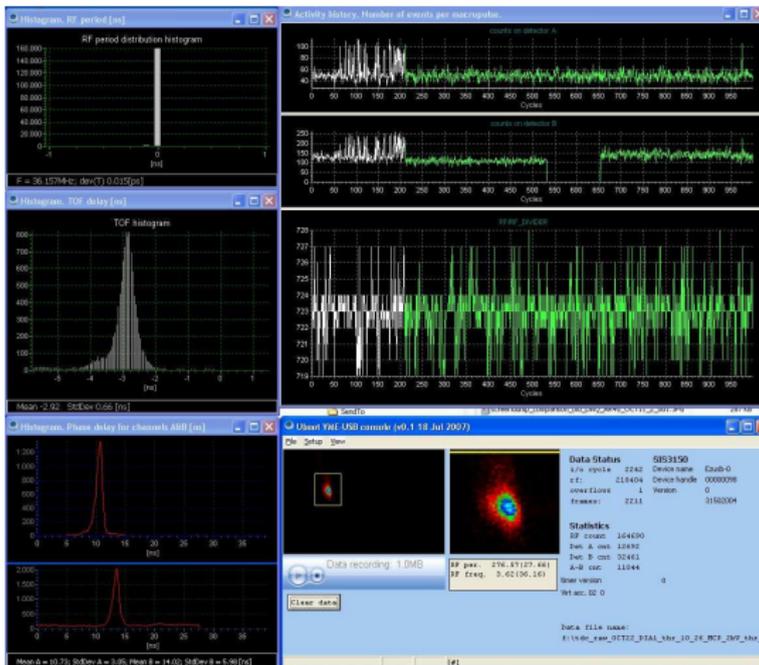
$\mu = 11334.9$  ticks  $\hat{=}$  10 rf periods

$\sigma \simeq 1.32$  ticks  $\hat{=}$  32.25 ps



Ticks during UNILAC rf periods (36.136 MHz, rf divider 10).

## DAQ GUI



# Primary Beam Attenuation Stage

Primary beam usually contains  $> 10^9$  particles per bunch.  
Beam attenuation is required due to

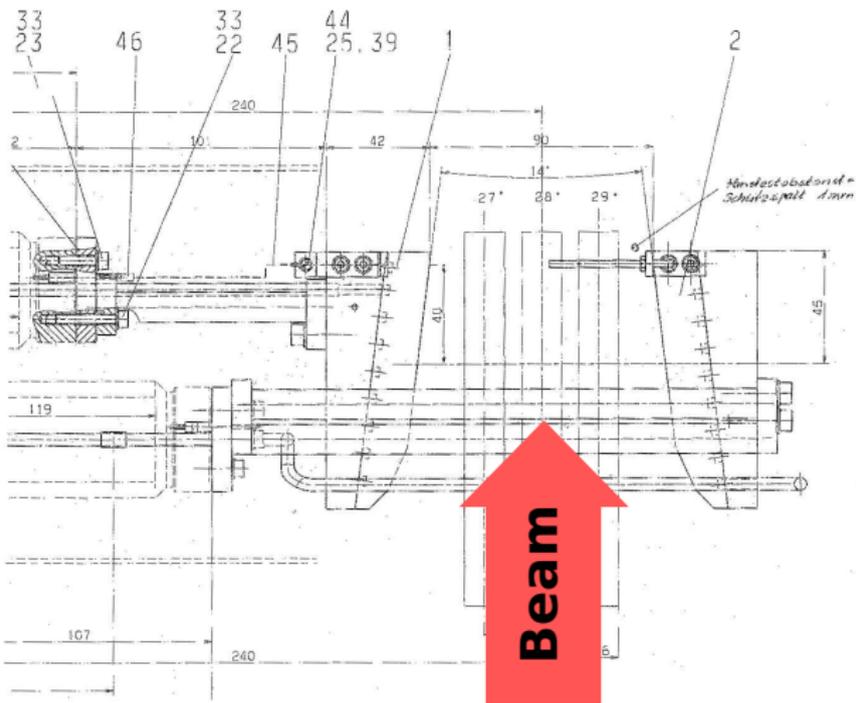
- sensitive foils (Ta  $210 \mu\text{g cm}^{-2}$ , Al  $9 \mu\text{g cm}^{-2}$ )
- required unambiguous coincidence condition, i.e. *rarely* more than one particle per micro pulse.

# Primary Beam Attenuation Stage

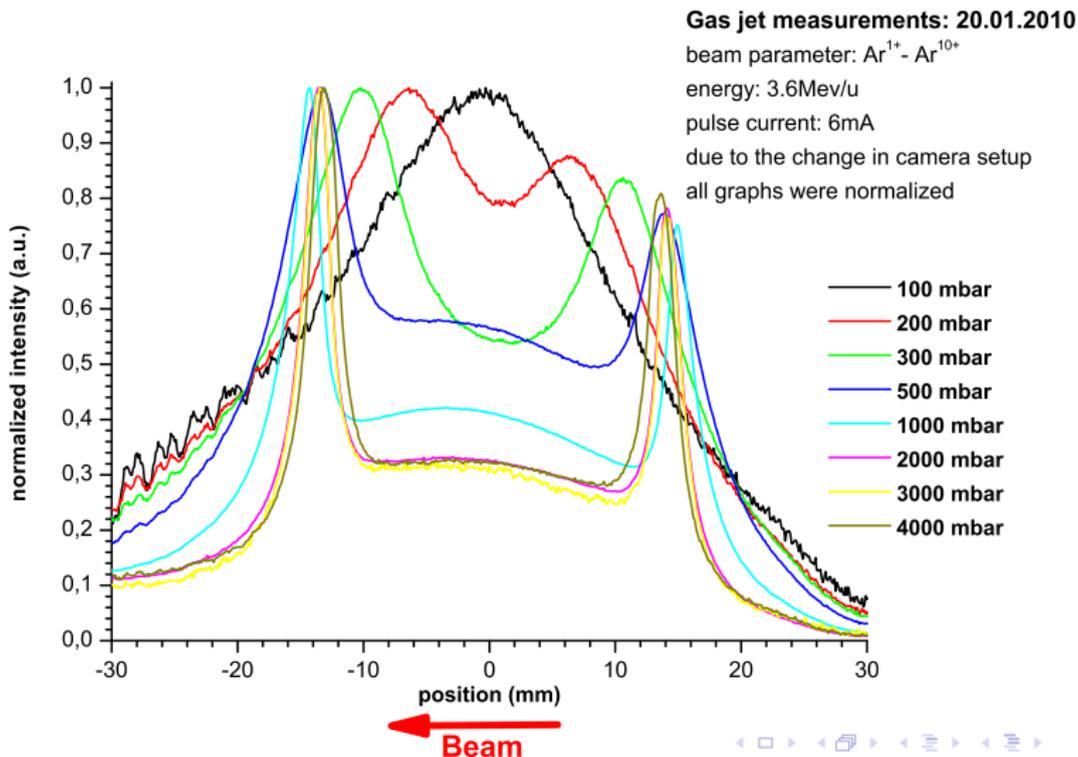
Current moderation down to several  $\mu A$  (20  $\mu A$  typical):

- Low/Zero current:
  - Defocusing of the primary beam by quadrupole magnet in front of the RFQ section.
  - Selection of a suitable charge state after stripping.
- High current (space charge effects):
  - Adjustment of the gas pressure inside the stripper section.
  - Selection of a suitable charge state after stripping.
  - High current slits acting as horizontal scraper.

# High Current Slits DS4, DS5 (horz. plane)



# Stripper: Gas Jet Profile (BIF)



# Secondary Beam Attenuation Stage: Particle attenuation via Coulomb scattering

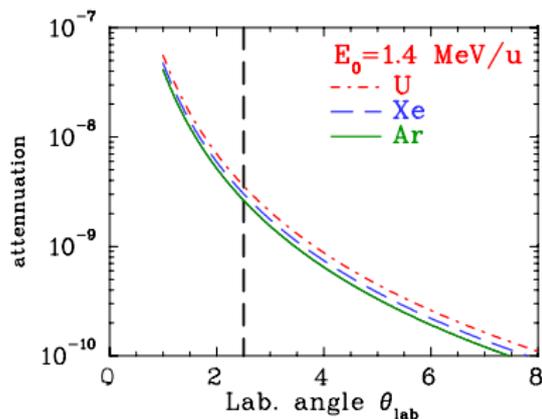
- Coulomb scattering via Ta foil
- Two collimators mounted to satisfy a scattering angle of  $2.5^\circ$ .

- Rutherford cross section:

$$\sigma_{Lab} = \left( \frac{Z_1 Z_2 e^2}{4E_C} \right)^2 \cdot \frac{1}{\sin^4 \frac{\theta_C}{2}}$$

- Attenuation:

$$\frac{N_{det}}{N_{bunch}} = \sigma_{Lab} \cdot d_{tar} \cdot \left( \frac{r_{col}}{r_{beam}} \right)^2 \cdot \delta\Omega$$



# Constructing the Phase Space from Time Stamps

2-dim histogramming using time stamps from MCP, Diamond and UNILAC rf.

- Spatial phase information from arrival time at Diamond with respect to UNILAC rf.

$$\Delta t_i = t_i^{DIA} - t_{rf}$$

- Energy/Momentum information from TOF between MCP and Diamond (80 cm).

$$\Delta E_i \approx \left. \frac{d}{dt} E(t) \right|_{t_{\langle E \rangle}} \cdot \Delta t_i = -2 \frac{\langle E \rangle}{t_{\langle E \rangle}} \cdot \Delta t_i^{TOF}$$

$$\Rightarrow \frac{\Delta E}{\langle E \rangle}_i \approx -2 \frac{\Delta t_i^{TOF}}{t_{\langle E \rangle}} = -\frac{2\beta c}{d_{TOF}} \cdot \left( t_i^{DIA} - t_i^{MCP} \right)$$

# Twiss Parameter - Short Overview

Integral statistic parameter to describe the beam.

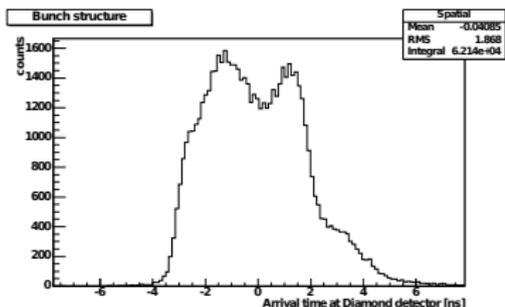
$$\varepsilon_{(rms)} = \sqrt{\langle x^2 \rangle \langle y^2 \rangle - \langle xy \rangle^2} \quad \text{RMS-Emittance}$$

$$\alpha = -\frac{\langle xy \rangle}{\varepsilon} = -\frac{\text{cov}(x, y)}{\varepsilon} \quad \text{Correlation}$$

$$\beta = \frac{\langle x^2 \rangle}{\varepsilon} = \frac{\sigma_x^2}{\varepsilon}$$

$$\gamma = \frac{\langle y^2 \rangle}{\varepsilon} = \frac{\sigma_y^2}{\varepsilon}$$

# Typical Measurement Example



file name: hippia1\_005.cor

$\langle E \rangle$ : 1.41 MeV/u  $\beta$ : 0.055

automatic Viewport boundaries  
 raw event data  
 integral count: 62960  
 bin size horizontal: 5 [LSB/TDC]  
 bin size vertical: 3 [LSB/TDC]

**Expected**

$\epsilon_{rms}$ : 42.96 keV ns

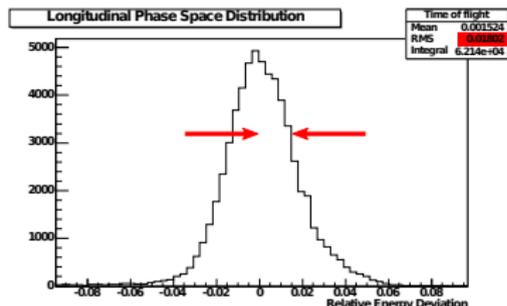
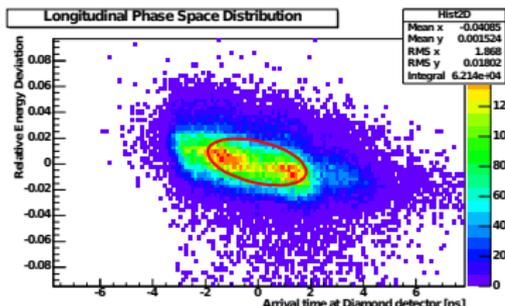
$\alpha$ : 0.467

$\beta$ : 0.0812 ns / keV

$\gamma$ : 15 keV / ns

$$\alpha \approx 4$$

$$\sigma_E \approx 0.0092$$

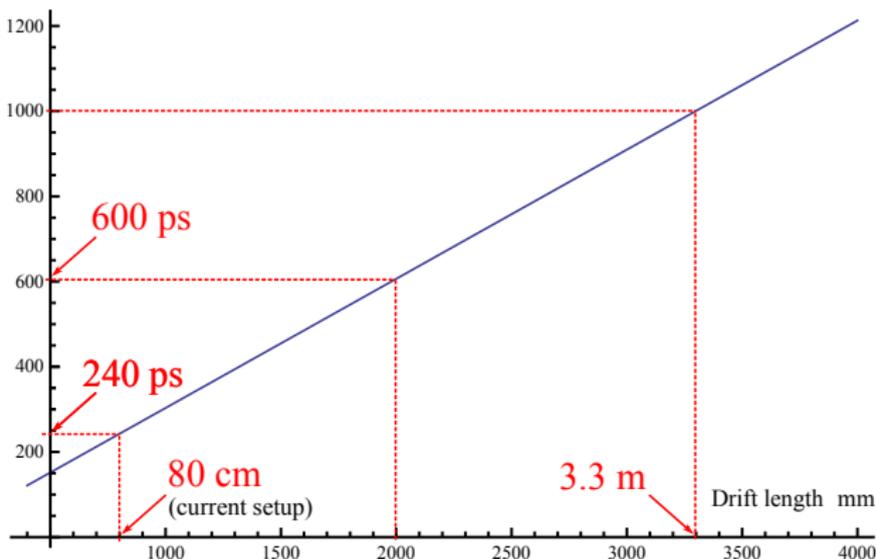


40AR 6.5 mA (60k events)

# Getting a glance of the required Time Resolution

*Difference* in TOF between a *design* particle at 1.4 MeV/u and a particle that deviates 1% (expected RMS) vs. detector separation.

TOF for rms energy spread [ps]



# Systematic Errors

The systematic error can be coarsely grouped into two categories.

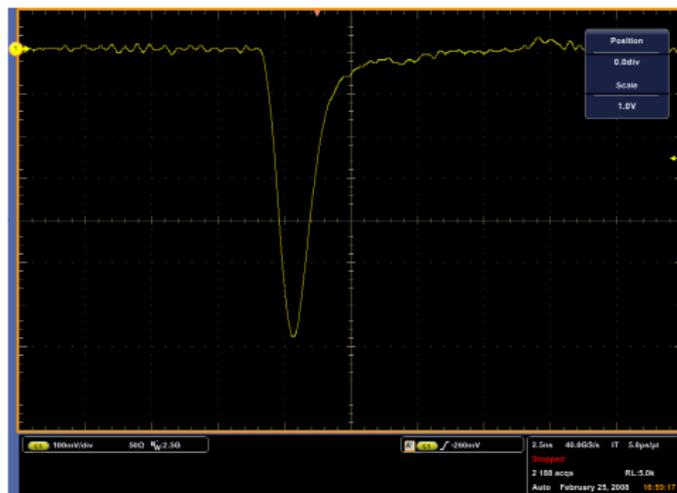
- Time resolution of detectors and electronic modules.
- *Physical* Phase Space distribution affected by the working principle of the device.

# Systematic Contributions from Detectors and Electronics

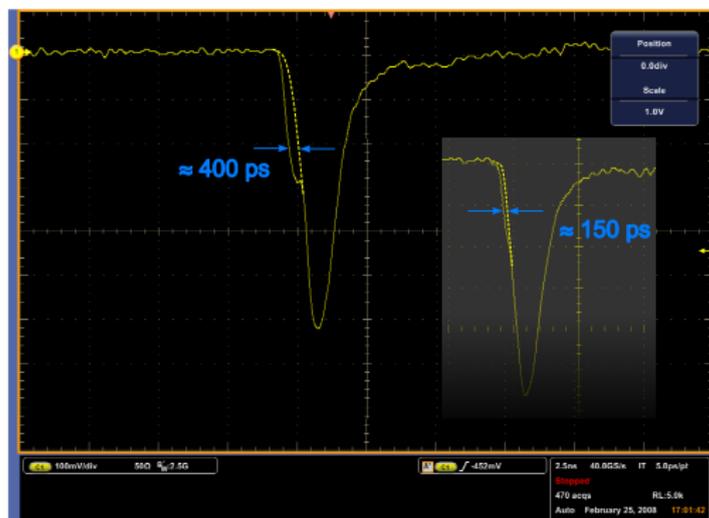
- TDC jitter per input link.
- Pulse shape discrimination jitter (walk, rise-time, edge distortions) from MCP and Diamond.
- Tilted Al-foil in front of MCP ( $37.5^\circ$ ) results in a spread in drift length which is used to determine the time-of-flight.
- Extension of Diamond electrode (8 mm) causes a jitter due to signal propagation depending on the location where a particle hits the electrode.
- Velocity distribution of emitted secondary  $e^-$  (MCP module).
- Anything major is missing?

# MCP signal shape - The Beauty

- Expected sane shape of MCP signal.
- Height varies between 300mV to 1.2V.

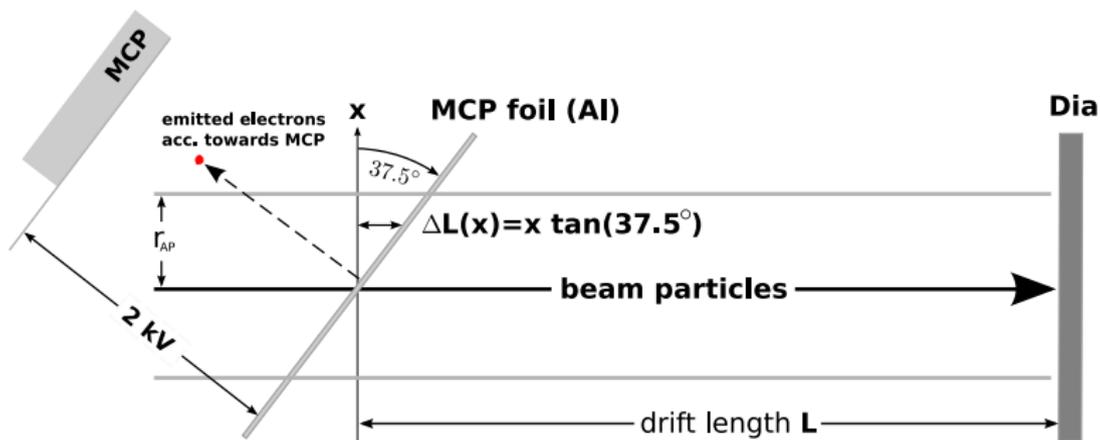


# MCP signal shape - And the Beast(s)



RMS jitter is difficult to estimate since no monochromatic beam is available with  $\frac{\Delta E}{\langle E \rangle} \ll 1\%$  that would allow a measurement of the pulse response.

# Tilted Al Foil inside MCP Module



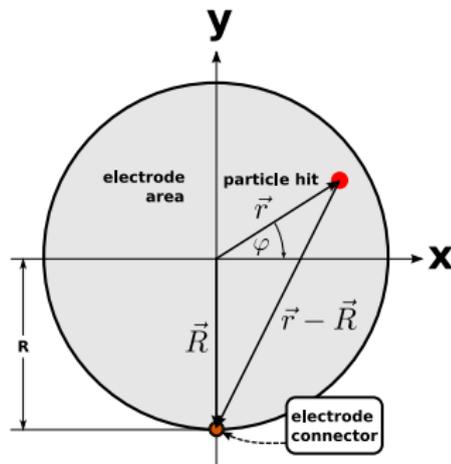
$\sigma_{\Delta L} = \frac{\tan\varphi}{\sqrt{3}} r_{AP}$ .  $r_{AP}$  is the effective aperture depending on the separation of the MCP and Diamond detector. For the current setup where the separation is 80 cm this contribution can be calculated as  $\approx 25$  ps.

# Signal Propagation on Diamond Electrode

$R=4$  mm, Signal propagation  $\approx c$

$$\sigma_{|\vec{r}-\vec{R}|} = \frac{\int_0^R r dr \int_0^{2\pi} d\varphi \left( \langle |\vec{r}-\vec{R}| \rangle - r \sqrt{\cos^2 \varphi + \left(\sin \varphi + \frac{R}{r}\right)^2} \right)^2}{\int_0^R r dr \int_0^{2\pi} d\varphi}$$

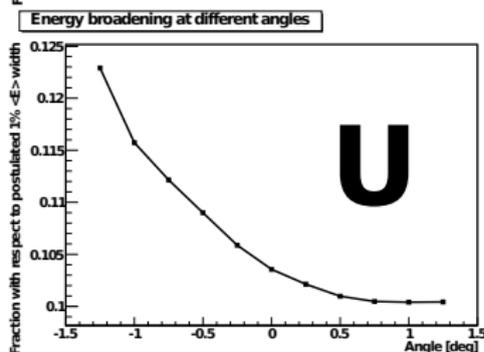
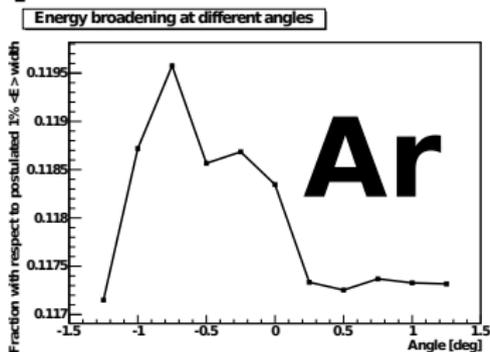
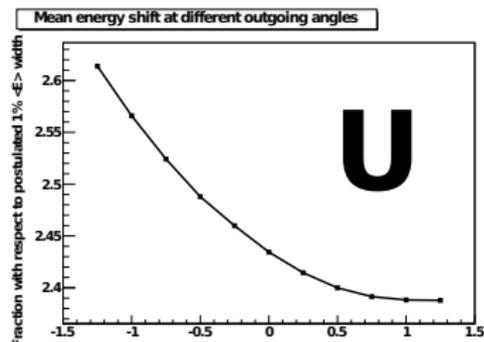
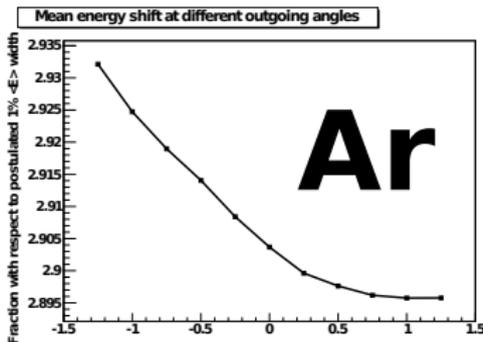
$$\Rightarrow t_{\sigma_{|\vec{r}-\vec{R}|}} = \frac{\sigma_{|\vec{r}-\vec{R}|}}{c} \approx 11 \text{ ps}$$



# Systematic Contributions from affected Phase Space

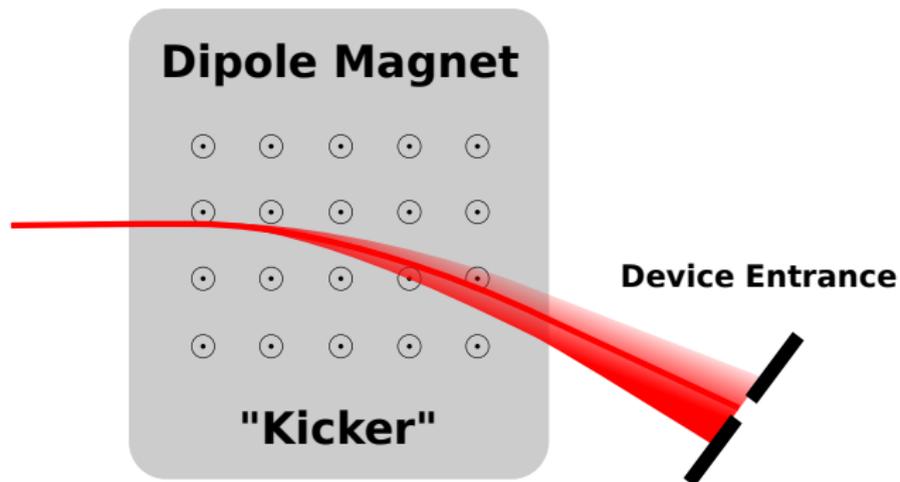
- Straggling contribution in Ta and Al foil (Electronic Stopping).
- Energy dependent scattering angle at Ta foil (finite solid angle).
- Attenuation via gas stripper and slits.
- T-L coupling in dispersive dipole section.
- Inhomogeneities in Ta and Al foil.
- Again, anything major is missing?

# Electronic Stopping SRIM calculations (Ar, U)

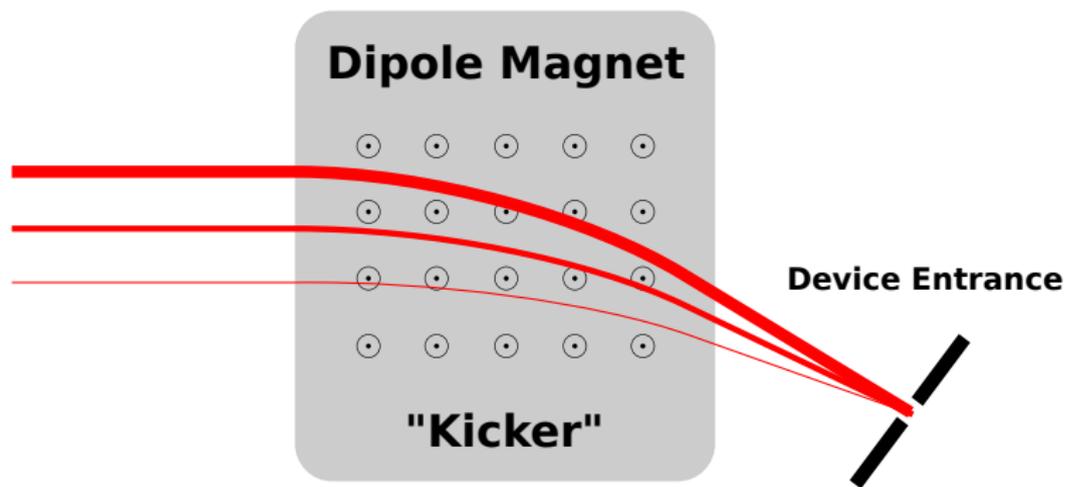


# Dispersive Section - Dipole ("Kicker")

The Lorentz Force...  $\vec{F}_L = q \cdot (\vec{E} + \vec{v} \times \vec{B})$



# Coupling of Transversal and Longitudinal Phase Space



# Coupling of Transversal and Longitudinal Phase Space

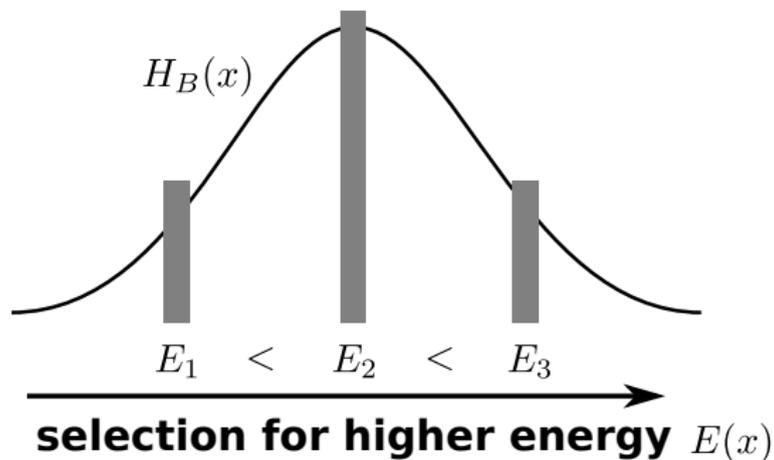
Can we treat the longitudinal and transversal Phase Space as decoupled degrees of freedom?

$$P(x, x', y, y', p, \varphi) \stackrel{?}{\approx} P(x, x', y, y') \otimes P(p, \varphi)$$

Even if we can - does that mean the dispersive section does *not* spoil the game?

# T-L Decoupled Phase Space

$H_B(x)$  : Transversal, horizontal beam profile

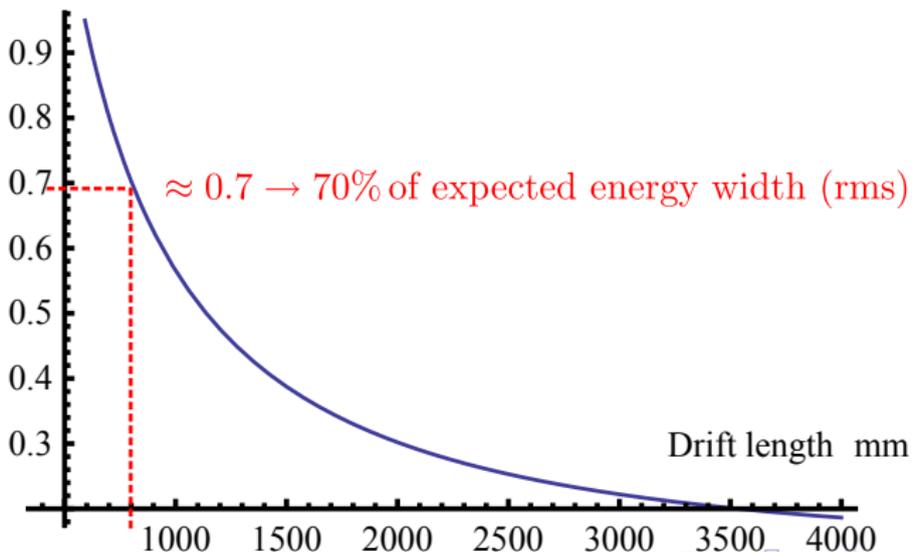


$$I_m(E) = H_B(x(E)) \cdot I_{Long}(E)$$

# Making an Estimate on Time Resolution from Systematic Effects

Taking into account all systematic effects that can be pinned down.

Error Calculated [% estimated rms energy width]



# Limited Resolution in a Gaussian Model Space

- Consider a 2-dim Gaussian phase space distribution given by Twiss parameters  $\{\alpha, \beta, \gamma, \varepsilon\}$

$$G(x, y; \alpha, \beta, \gamma, \varepsilon) = \frac{1}{2\pi\varepsilon} e^{-\frac{1}{2\varepsilon}(\gamma x^2 + 2\alpha xy + \beta y^2)}$$

- Assume a systematic *Gaussian pulse response* (for now  $y$  only)

$$g(y; \sigma_{res}) = \frac{1}{\sqrt{2\pi}\sigma_{res}} e^{-\frac{y^2}{2\sigma_{res}^2}}$$

# Convolved Phase Space Distribution

- Allows qualitative analytic investigations on how *measured* Twiss parameters are affected by means of a convolution

$$\begin{aligned}
 (g * G)(x, y) &= \int_{-\infty}^{+\infty} dz G(x, z; \alpha, \beta, \gamma, \varepsilon) \cdot g(y - z; \sigma_{res}) \\
 &= \frac{1}{2\pi\varepsilon\kappa} e^{-\frac{1}{2\varepsilon\kappa} \left\{ \left( \gamma\kappa - \frac{\sigma_{res}^2}{\varepsilon} \frac{\alpha^2}{\kappa} \right) x^2 + 2\frac{\alpha}{\kappa} xy + \frac{\beta}{\kappa} y^2 \right\}}
 \end{aligned}$$

$$\text{with } \kappa := \sqrt{\sigma_{res}^2 \frac{\beta}{\varepsilon} + 1} \quad \Rightarrow \quad (\kappa > 1) \text{ for } \sigma_{res} \neq 0$$

# Transformation of Twiss parameters

- Invariant *form* of representation allows comparison of coefficients with undisturbed phase space and thus provides the transformation of the Twiss parameters depending on  $\sigma_{res}$ .

$$\varepsilon' = \varepsilon \cdot \sqrt{\sigma_{res}^2 \frac{\beta}{\varepsilon} + 1} = \varepsilon \cdot \kappa$$

$$\alpha' = \frac{\alpha}{\kappa} = \alpha \cdot \frac{\varepsilon}{\varepsilon'}$$

$$\beta' = \frac{\beta}{\kappa} = \beta \cdot \frac{\varepsilon}{\varepsilon'}$$

$$\gamma' = \gamma \cdot \kappa - \frac{\sigma_{res}^2}{\varepsilon} \cdot \frac{\alpha^2}{\kappa} = \gamma \cdot \frac{1 + \frac{\kappa^2 - 1}{\alpha^2 + 1}}{\kappa}$$

$$\Rightarrow \beta' \cdot \gamma' - \alpha'^2 = 1 \quad \checkmark$$

# Transformation of Twiss parameters

- General tendency of Twiss parameters in case of vertical resolution limit.

$$\varepsilon' > \varepsilon$$

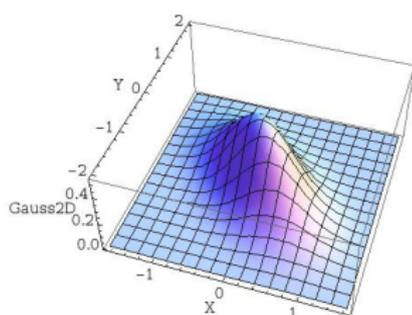
$$\alpha' < \alpha$$

$$\beta' < \beta$$

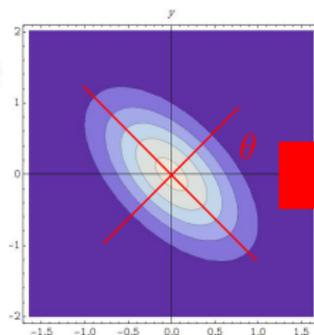
$$\gamma' = \begin{cases} \leq \gamma, & \kappa \leq \alpha^2 \text{ and } \alpha > 1 \\ > \gamma & \text{otherwise} \end{cases}$$

# 2-dim Gaussian Sample Distribution with Vertical Gaussian Resolution Limit $\sigma_{res} = \sigma_y$

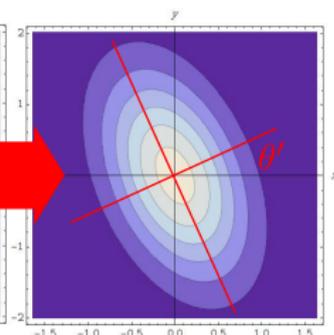
Original:  $\varepsilon = 0.3$ ,  $\alpha = 0.71$ ,  $\beta = 1$ ,  $\gamma = 1.5$ ,  $\sigma_y = 0.67$ ,  $\theta = 35.3^\circ$



Original



Original



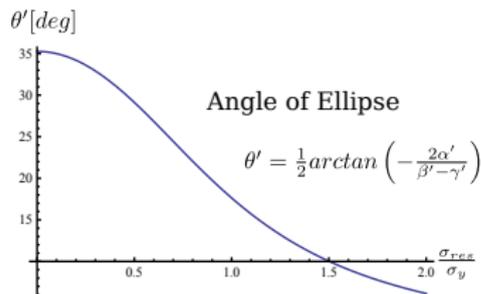
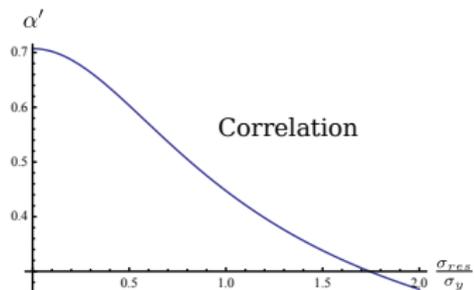
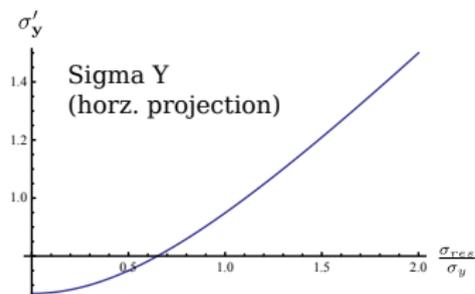
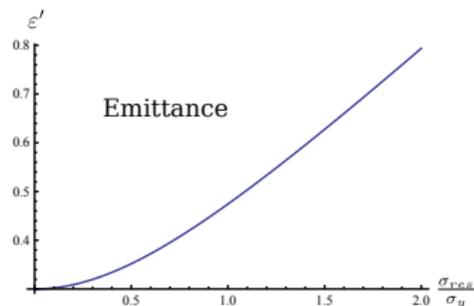
Convolved

Parameters after convolution:

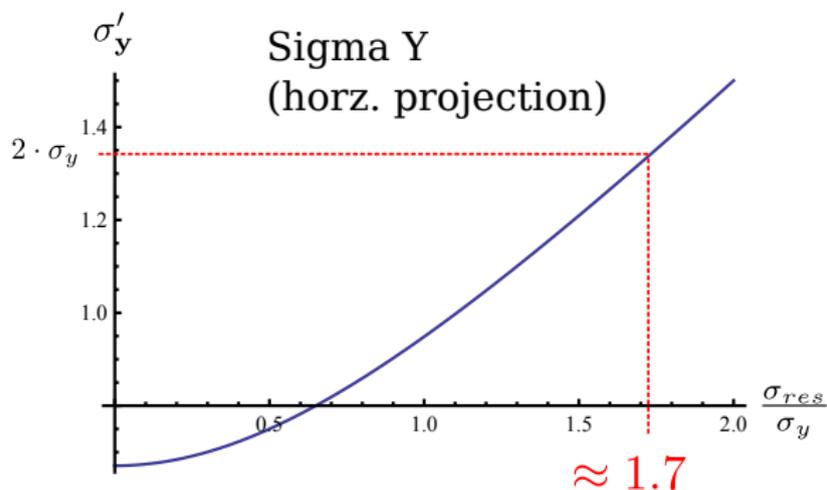
$\varepsilon' = 0.48$ ,  $\alpha' = 0.45$ ,  $\beta' = 0.63$ ,  $\gamma' = 1.9$ ,  $\sigma_y' = 0.95$ ,  $\theta' = 17.6^\circ$

# Beam Parameters vs. Vertical Resolution

Beam parameter:  $\varepsilon = 0.3$ ,  $\beta = 1$ ,  $\gamma = 1.5$  ( $\rightarrow \alpha = \frac{1}{\sqrt{2}}$ )



# Estimation of Vertical Resolution



- Compared to our previous estimation ( $\approx 0.7$ ) we might be missing something important.
- (Or) Do we underestimate some contribution?
- E.g. variation in thickness of primary scatter foil.

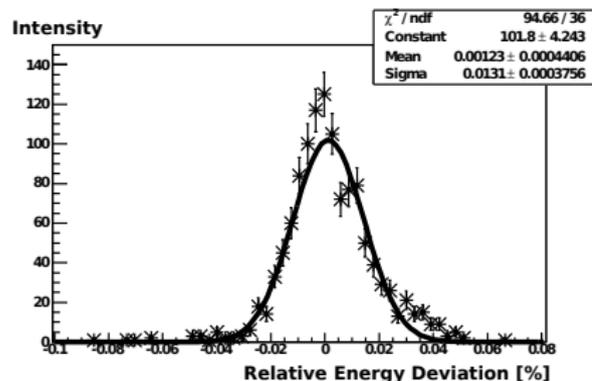
# How about “common” Deconvolution then?

- Deconvolution is the attempt to reverse a convolution between to functions.
- Requires the precise knowledge of the pulse response (resolution property).
- *Very* sensitive to SNR.
- Difficult to find a consistent method to fit our needs with low statistics.

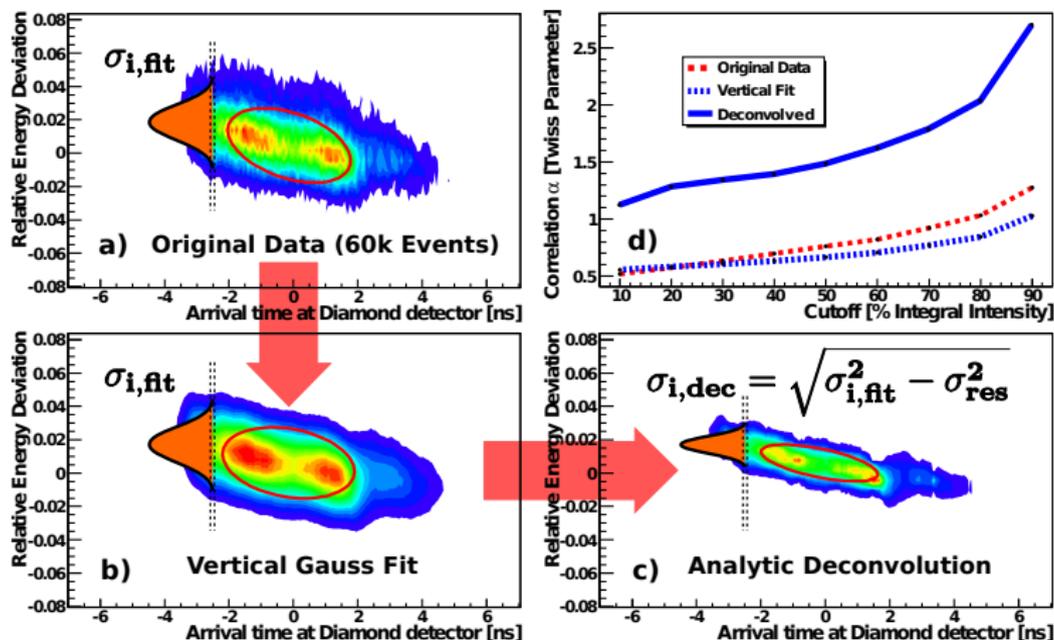
# Analytic Deconvolution

- Fit vertical energy slice 'i' to a Gaussian parameterization ( $\sigma_{i,fit}^2$ ).
- Pulse response  $\sigma_{res}$  deduced from error estimation is also considered Gaussian due to the different systematic contributions.
- Analytic deconvolution via Gaussian convolution theorem  

$$\sigma_{i,dec} = \sqrt{\sigma_{i,fit}^2 - \sigma_{res}^2}$$
- Reconstruct Phase Space with  $\{\sigma_{i,dec}\}$

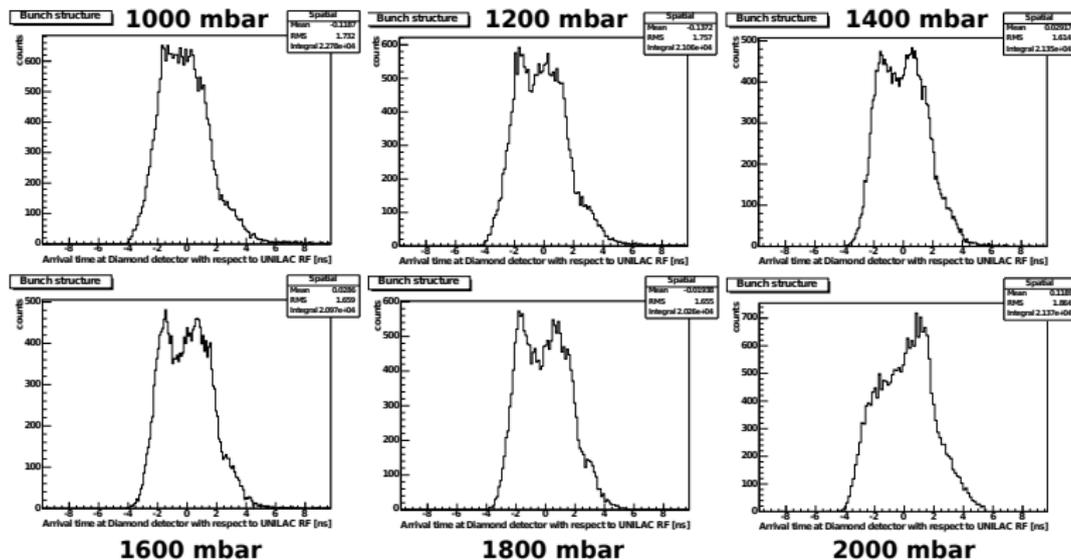


# Analytic Deconvolution



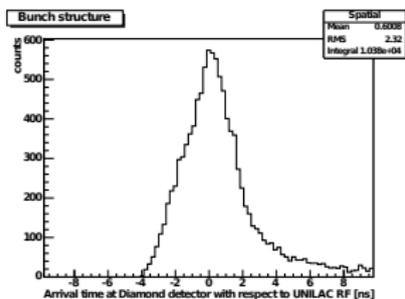
# Variation of Stripper Gas Pressure

40Ar, high current setup, DS4 +/- 15 mm, DS5 +/- 6mm  
constant



## 500 mbar

40Ar, high current setup, DS4 +/- 15 mm, DS5 +/- 6mm  
constant



file name: AR14p\_300mus\_470mbar.cor

$\langle E \rangle$ : 1.41 MeV/u  $\beta$ : 0.055

automatic Viewport boundaries

raw event data

integral count: 11240

bin size horizontal: 9 [LSB/TDC]

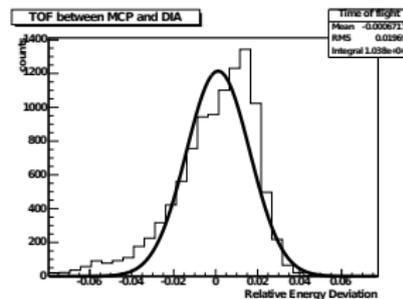
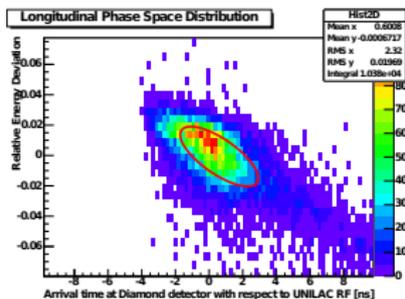
bin size vertical: 5 [LSB/TDC]

$\epsilon_{rms}$ : 44.20 keV ns

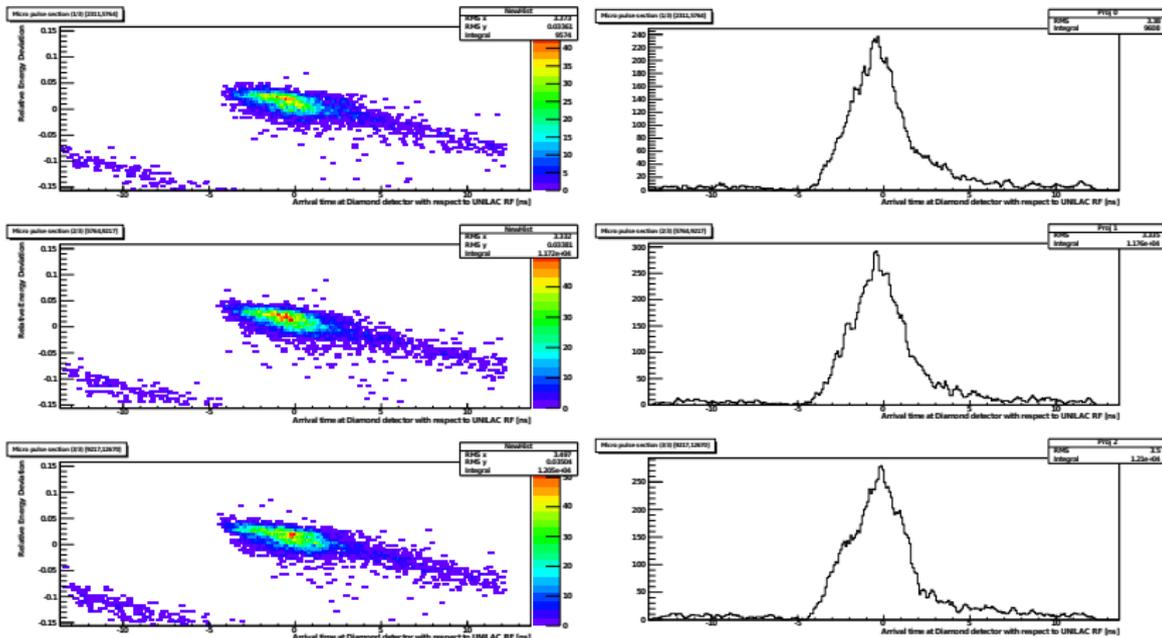
$\alpha$ : 1.06

$\beta$ : 0.122 ns / keV

$\gamma$ : 17.4 keV / ns

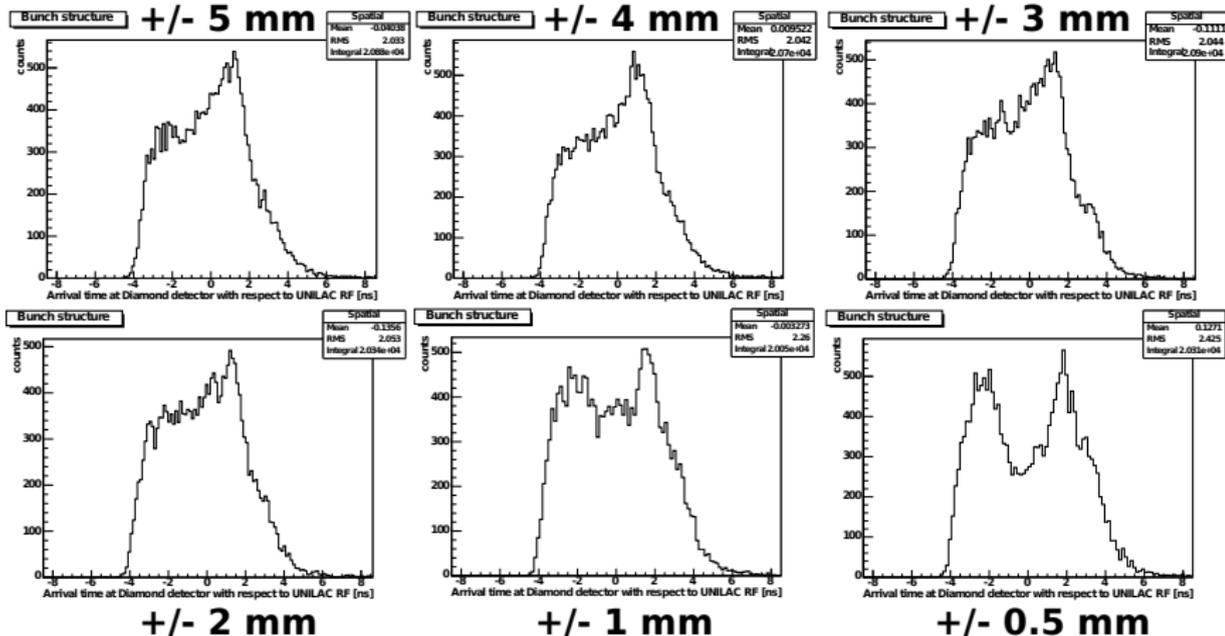


# Micro Pulse Cuts (500 mbar)



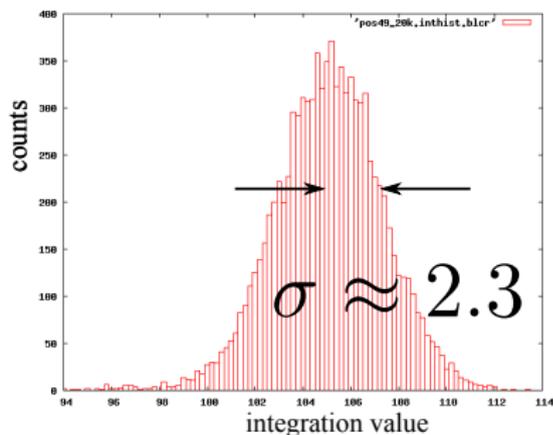
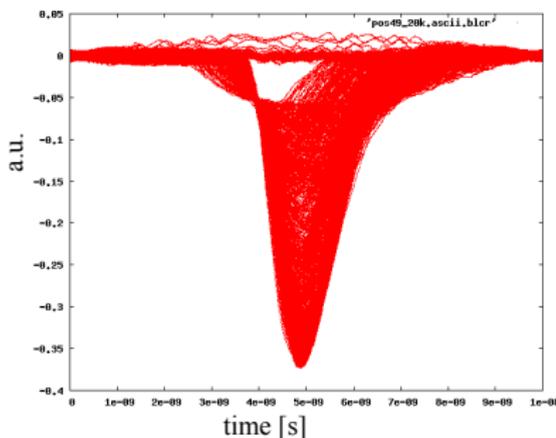
# Variation of Slit DS5 (DS4 +/- 15 mm const.)

40Ar, high current setup, gas pressure 2000 mbar (constant)



# Preliminary: Direct Calorimetric Measurement

Am-243 Source, vac  $10^{-5}$  mbar



Tektronix DPO 7542 - 2.5 GHz BW, 20 GS/s (2 ch.)

# Pre-Preliminary: Discrimination via Post-Processing

- Sample full data stream with Tek 7254 and discriminate afterwards.
- Can we perform better with this than the double threshold discriminators do yet?
- If so - does it improve the overall situation significantly?

# Outlook



**Thanks for your attention!**  
**Questions?**

