# Beam Position Monitor: <br> Detector Principle, Hardware and Electronics Peter Forck, Piotr Kowina and Dmitry Liakin Gesellschaft für Schwerionenforschung, Darmstadt 

## Outline:

$>$ Signal generation $\rightarrow$ transfer impedance
$>$ Consideration for capacitive shoe box BPM
$>$ Consideration for capacitive button BPM
$>$ Other BPM principles: stripline $\rightarrow$ traveling wave inductive $\rightarrow$ wall current cavity $\rightarrow$ resonator for dipole mode
$>$ Electronics for position evaluation
$>$ Some examples for position evaluation and other applications
$>$ Summary

## General Idea: Detection of Wall Charges

The image current at the vacuum wall is monitored on a high frequency basis i.e. the ac-part given by the bunched beam.


For relativistic velocities, the electric field is mainly transversal: $E_{\perp, l a b}(t)=\gamma \cdot E_{\perp, \text { rest }}(t)$


## Usage of BPMs

## A BPM is an non-destructive device

It has a low cut-off frequency i.e. dc-beam behavior can not be monitored (exception: Schottky spectra, here the physics is due to finite number of particles)
$\Rightarrow$ Usage with bunched beams!
It delivers information about:

1. The center of the beam
> Closed orbit
i.e. central orbit averaged over a period much longer than a betatron oscillation
$>$ Bunch position on a large time scale: bunch-by-bunch $\rightarrow$ turn-by-turn $\rightarrow$ averaged position
$>$ Single bunch position $\rightarrow$ determination of parameters like tune, chromaticity, $\beta$-function
$>$ Time evolution of a single bunch can be compared to 'macro-particle tracking' calculations
$>$ Feedback: fast bunch-by-bunch damping $\rightarrow$ precise (and slow) closed orbit correction

## 2. Longitudinal bunch shapes

$>$ Bunch behavior during storage and acceleration
$>$ For proton LINACs: the beam velocity can be determined by two BPMs
> Relative low current measurement down to 10 nA .

## Model for Signal Treatment of capacitive BPMs

The wall current is monitored by a plate or ring inserted in the beam pipe:


The image current $\mathrm{I}_{\mathrm{im}}$ at the plate is given by the beam current and geometry:

$$
\begin{aligned}
I_{i m}(t) & \equiv d Q_{i m} / d t=A / 2 \pi a l \cdot d Q_{\text {beam }}(t) / d t=A / 2 \pi a l \cdot l / \beta c \cdot d I_{\text {beam }} / d t \\
& =A / 2 \pi a \cdot 1 / \beta c \cdot i \omega I_{\text {beam }}(\omega)
\end{aligned}
$$

Using a relation for Fourier transformation: $I_{\text {beam }}=I_{0} e^{i \omega t} \Rightarrow d I_{\text {beam }} d t=i \omega I_{\text {beam }}$.

## Transfer Impedance for capacitive BPM

At a resistor R the voltage $U_{i m}$ from the image current is measured.
The transfer impedance $Z_{t}$ is the ratio between voltage $U_{i m}$ and beam current $I_{\text {beam }}$
in frequency domain: $U_{i m}(\omega)=R \cdot I_{i m}(\omega)=Z_{t}(\omega, \beta) \cdot I_{\text {beam }}(\omega)$.

## Capacitive BPM:

equivalent circuit
-The pick-up capacitance $C$ :
plate $\leftrightarrow$ vacuum-pipe and cable.
-The amplifier with input resistor $R$.
-The beam is a high-impedance current source:

$$
\begin{aligned}
U_{\text {im }} & =\frac{R}{1+i \omega R C} \cdot I_{i m} \\
& =\frac{A}{2 \pi a} \cdot \frac{1}{\beta c} \cdot \frac{1}{C} \cdot \frac{i \omega R C}{1+i \omega R C} \cdot I_{\text {beam }} \\
& \equiv Z_{t}(\omega, \beta) \cdot I_{\text {beam }}
\end{aligned}
$$



$$
\frac{1}{Z}=\frac{1}{R}+i \omega C \Leftrightarrow Z=\frac{R}{1+i \omega R C}
$$

This is a high-pass characteristic with $\omega_{\text {cut }}=1 / R C$ :
Amplitude: $\left|Z_{t}(\omega)\right|=\frac{A}{2 \pi a} \cdot \frac{1}{\beta c} \cdot \frac{1}{C} \cdot \frac{\omega / \omega_{\text {cut }}}{\sqrt{1+\omega^{2} / \omega_{c u t}^{2}}}$ Phase: $\varphi(\omega)=\arctan \left(\omega_{c u t} / \omega\right)$

## Example of Transfer Impedance for Proton Synchrotron

The high-pass characteristic for typical synchrotron BPM:

$$
\begin{aligned}
& U_{\text {im }}(\omega)=Z_{t}(\omega) \cdot I_{\text {beam }}(\omega) \\
& \left|Z_{t}\right|=\frac{A}{2 \pi a} \cdot \frac{1}{\beta c} \cdot \frac{1}{C} \cdot \frac{\omega / \omega_{\text {cut }}}{\sqrt{1+\omega^{2} / \omega_{c u t}^{2}}} \\
& \varphi=\arctan \left(\omega_{c u t} / \omega\right)
\end{aligned}
$$

Parameter for shoe-box BPM:

$$
C=100 \mathrm{pF}, l=10 \mathrm{~cm}, \beta=50 \%
$$

$$
f_{\text {cut }}=\omega / 2 \pi=(2 \pi R C)^{-1}
$$

$$
\text { for } R=50 \Omega \Rightarrow f_{\text {cut }}=32 \mathrm{MHz}
$$

$$
\text { for } R=1 \mathrm{M} \Omega \Rightarrow f_{\text {cut }}=1.6 \mathrm{kHz}
$$



Large signal strength $\rightarrow$ high impedance
Smooth signal transmission $\boldsymbol{\rightarrow 5 0} \Omega$
Compromise $\rightarrow \approx 5 \mathrm{k} \Omega$ by transformer e.g. $N_{\text {prim }} / N_{\text {sec }}=3: 30$
Impedance $Z_{\text {prim }}=\left(N_{\text {prim }} / N_{\text {sec }}\right)^{2} \cdot Z_{\text {sec }}$ voltage $U_{\text {im }}=N_{\text {sec }} / N_{\text {prim }} \cdot U_{\text {prim }}$
$\rightarrow$ Smooth signal chain, medium cut-off frequency, but lower usable voltage


## Transfer Impedance Measurement

With a network analyzer and an antenna the BPM properties can be determined.


Signal Shape for capacitive BPMs: differentiated $\leftrightarrow$ proportional

Depending on the frequency range and termination the signal looks different:
$>$ High frequency range $\omega \gg \omega_{\text {cut }}{ }^{\circ}$

$$
Z_{t} \propto \frac{i \omega / \omega_{c u t}}{1+i \omega / \omega_{c u t}} \rightarrow 1 \Rightarrow U_{i m}(t)=\frac{1}{C} \cdot \frac{1}{\beta c} \cdot \frac{A}{2 \pi a} \cdot I_{\text {beam }}(t)
$$

$\Rightarrow$ direct image of the bunch. Signal strength $Z_{t} \alpha A / C$ i.e. nearly independent on length
$>$ Low frequency range $\omega \ll \omega_{\text {cut }}$.
$Z_{t} \propto \frac{i \omega / \omega_{\text {cut }}}{1+i \omega / \omega_{\text {cut }}} \rightarrow i \frac{\omega}{\omega_{\text {cut }}} \Rightarrow U_{\text {im }}(t)=R \cdot \frac{A}{\beta c \cdot 2 \pi a} \cdot i \omega I_{\text {beam }}(t)=R \cdot \frac{A}{\beta c \cdot 2 \pi a} \cdot \frac{d I_{\text {beam }}}{d t}$
$\Rightarrow$ derivative of bunch, single strength $\mathrm{Z}_{\mathrm{t}} \alpha \mathrm{A}$, i.e. (nearly) independent on C
Example from synchrotron BPM with $50 \Omega$ termination (reality at p-synchrotron : $\sigma \gg 1 \mathrm{~ns}$ ):


## Examples for differentiated \& proportional Shape

Proton LINAC, $\mathrm{e}^{-}-$LINAC\&synchtrotron: $100 \mathrm{MHz}<f_{r f}<1 \mathrm{GHz}$ typically $R=50 \Omega$ processing to reach bandwidth $C \approx 5 \mathrm{pF} \Rightarrow f_{\text {cut }}=1 /(2 \pi R C) \approx 700 \mathrm{MHz}$ Example: 36 MHz GSI ion LINAC


Proton synchtrotron:
$1 \mathrm{MHz}<f_{r f}<30 \mathrm{MHz}$ typically $R=1 \mathrm{M} \Omega$ for large signal i.e. large $\mathrm{Z}_{\mathrm{t}}$ $C \approx 100 \mathrm{pF} \Rightarrow f_{\text {cut }}=1 /(2 \pi R C) \approx 10 \mathrm{kHz}$ Example: non-relativistic GSI synchrotron $f_{r f}: 0.8 \mathrm{MHz} \rightarrow 5 \mathrm{MHz}$



Remark: During acceleration the bunching-factor is increased: 'adiabatic damping'.

## Example Shoe-box BPMs

Shoe-box BPMs used at low $\beta$ proton \& ion synchrotron for $1 \mathrm{MHz}<f_{r f}<10 \mathrm{MHz}$.
Example: HIT cancer therapy synchrotron $0.8 \mathrm{MHz}<f_{r f}<5 \mathrm{MHz}$


## Calculation of Signal Shape: Single Bunch

The transfer impedance is used in frequency domain! The following is performed:

1. Start: Time domain Gaussian function $I_{\text {beam }}(t)$ having a width of $\sigma_{t}$

2. FFT of $I_{\text {beam }}(t)$ leads to the frequency domain Gaussian $I_{\text {beam }}(f)$ with $\sigma_{f}=\left(2 \pi \sigma_{t}\right)^{-1}$
$I_{\text {beam }}(f)$
$U_{\text {im }}(f)$

3. Multiplication with $Z_{t}(f)$ with $f_{\text {cut }}=32 \mathrm{MHz}$ leads to $U_{\text {im }}(f)=Z_{t}(f) \cdot I_{\text {beam }}(f)$
4. Inverse FFT leads to $U_{\text {im }}(t)$

## Calculation of Signal Shape: Bunch Train

Example for low energy proton synchrotron: Train of bunches with $\mathrm{R}=1 \mathrm{M} \Omega$




Calculation: $I_{\text {beam }}(t) \xrightarrow{\text { FFT }} I_{\text {beam }}(\omega) \rightarrow U_{\text {im }}(\omega)=Z_{\text {tot }}(\omega) \cdot I_{\text {beam }}(\omega) \xrightarrow{\text { invFFT }} U_{\text {im }}(t)$
Parameter: $R=1 \mathrm{M} \Omega \Rightarrow f_{\text {cut }}=2 \mathrm{kHz}, Z_{t}=5 \Omega$ all buckets filled, no amp

$$
C=100 \mathrm{pF}, l=10 \mathrm{~cm}, \beta=50 \%, \sigma_{t}=100 \mathrm{~ns} \Rightarrow \sigma_{l}=15 \mathrm{~m}
$$

$>$ Fourier spectrum is composed of lines separated by acceleration $f_{r f}$
$>$ Envelope given by single bunch Fourier transformation
$>$ Baseline shift due to ac-coupling
Remark: $1 \mathrm{MHz}<f_{r f}<10 \mathrm{MHz} \Rightarrow$ Bandwidth $\approx 100 \mathrm{MHz}=10 \cdot f_{r f}$ for broadband observation

## Calculation of Signal Shape: Bunch Train

Train of bunches with $\mathrm{R}=50 \Omega$ termination:




Parameter: $R=50 \Omega \Rightarrow f_{\text {cut }}=\mathbf{3 2} \mathrm{MHz}$, all buckets filled, no amp
$C=100 \mathrm{pF}, l=10 \mathrm{~cm}, \beta=50 \%, \sigma_{t}=100 \mathrm{~ns}$
$>$ Low frequency cut-off due to $f_{\text {cut }}=32 \mathrm{MHz}$
$>$ Differentiated bunches, 15 fold lower amplitude
$>$ Modified Fourier spectrum with low amplitude value, maximum shift to higher frequencies

## Calculation of Signal Shape: Bunch Train with empty Buckets

Synchrotron during filling: Empty buckets, $\mathrm{R}=5 \mathrm{k} \Omega$ termination:




Parameter: $R=5 \mathrm{k} \Omega \Rightarrow f_{\text {cut }}=320 \mathrm{kHz}, 2$ empty buckets
$C=100 \mathrm{pF}, l=10 \mathrm{~cm}, \beta=50 \%, \sigma=100 \mathrm{~ns}$
$>$ Fourier spectrum is more complex, harmonics are broader
$>$ Varying baseline with $\tau \approx\left(3 f_{\text {cut }}\right)^{-1}=1 \mu \mathrm{~s}$
$>$ Baseline shift calls for dedicated restoring algorithm for time domain processing.

## Calculation of Signal Shape: Bunch Train with Cable Damping

Effect of cable or other electronics:



$>$ Bunch signal is damped; 10 fold lower amplitude, higher frequencies are damped stronger
$>$ Bunch signal gets asymmetric, baseline did not reach zero
$>\Rightarrow$ 'Good cables' are a precaution for broadband signal transmission

## Calculation of Signal Shape: Filtering of Harmonics

Effect of filters, here bandpass:




Parameter: $R=5 \mathrm{k} \Omega, 4{ }^{\text {III }}$ order Butterworth filter at $f_{\text {cut }}=\mathbf{2 ~ M H z}$

$$
C=100 \mathrm{pF}, l=10 \mathrm{~cm}, \beta=50 \%, \sigma=100 \mathrm{~ns}
$$

$>$ Ringing due to sharp cutoff
$>$ Other filter types more appropriate

$$
\begin{aligned}
& n^{\text {th }} \text { order Butterworth filter, math. simple, but not well suited: } \\
& \left|H_{\text {low }}\right|=\frac{1}{\sqrt{1+\left(\omega / \omega_{\text {cut }}\right)^{2 n}}} \text { and }\left|H_{\text {high }}\right|=\frac{\left(\omega / \omega_{\text {cut }}\right)^{n}}{\sqrt{1+\left(\omega / \omega_{\text {cut }}\right)^{2 n}}} \\
& H_{\text {filter }}=H_{\text {high }} \cdot H_{\text {low }}
\end{aligned}
$$

Generally: $\quad Z_{\text {tot }}(\omega)=H_{\text {cable }}(\omega) \cdot H_{\text {filter }}(\omega) \cdot H_{\text {amp }}(\omega) \cdot \ldots \cdot Z_{t}(\omega)$
Remark: For electronics calculation, time domain filters (FIR and IIR) are more appropriate

## Principle of Position Determination with BPM

The difference between plates gives the beam's center-of-mass

## $\rightarrow$ most frequent application

'Proximity' effect leads to different voltages at the plates:

$$
\begin{aligned}
& y=\frac{1}{S_{y}(\omega)} \cdot \frac{U_{u p}-U_{\text {down }}}{U_{u p}+U_{\text {down }}}+\delta_{y}(\omega) \\
& \equiv \frac{1}{S_{y}} \cdot \frac{\Delta U_{y}}{\Sigma U_{y}}+\delta_{y} \\
& x=\frac{1}{S_{x}(\omega)} \cdot \frac{U_{\text {pick up }}}{U_{\text {right }}-U_{\text {left }}}+\delta_{x}(\omega) \\
& U_{\text {right }}+U_{\text {left }}
\end{aligned}
$$

$\boldsymbol{S}(\boldsymbol{\omega}, \boldsymbol{x})$ is called position sensitivity, sometimes the inverse is used $\boldsymbol{k}(\boldsymbol{\omega}, \boldsymbol{x})=\mathbf{1 / S}(\boldsymbol{\omega}, \boldsymbol{x})$
$\boldsymbol{S}$ is a geometry dependent, non-linear function, which have to be optimized.
Units: $\boldsymbol{S}=[\% / \mathrm{mm}]$ and sometimes $\boldsymbol{S}=[\mathrm{dB} / \mathrm{mm}]$ or $\boldsymbol{k}=[\mathrm{mm}]$
Sometimes the transverse transfer impedance is defined via $\boldsymbol{U} \boldsymbol{\Delta}=\boldsymbol{Z}_{\perp}(\boldsymbol{\omega}) \cdot \boldsymbol{x} \boldsymbol{I}_{\text {beam }}$
It can be assumed: $Z_{\perp}(\omega, x)=Z_{t}(\omega) / \boldsymbol{S}(\omega, x)$

# Beam Position Monitor: <br> Detector Principle, Hardware and Electronics 

Outline:
$>$ Signal generation $\rightarrow$ transfer impedance
$>$ Consideration for capacitive 'shoe box' = 'linear cut' BPM position sensitivity calculation, crosstalk, realization
$>$ Consideration for capacitive button BPM
$>$ Other BPM principles: stripline $\rightarrow$ traveling wave
inductive $\rightarrow$ wall current
cavity $\rightarrow$ resonator for dipole mode
$>$ Electronics for position evaluation
$>$ Some examples for position evaluation and other applications
$>$ Summary

## Shoe-box BPM for Proton or Ion Synchrotron

Frequency range: $1 \mathrm{MHz}<\mathrm{f}_{\mathrm{rf}}<10 \mathrm{MHz} \Rightarrow$ bunch-length $\gg$ BPM length.

Signal is proportional to actual plate length:

$$
\begin{aligned}
& l_{\mathrm{right}}=(a+x) \cdot \tan \alpha \\
& \Rightarrow x=a \cdot \frac{l_{\mathrm{right}}-l_{\mathrm{left}}}{l_{\mathrm{right}}+l_{\mathrm{left}}}
\end{aligned}
$$

In ideal case: linear reading
$x=a \cdot \frac{U_{\text {right }}-U_{\text {left }}}{U_{\text {right }}+U_{\text {left }}} \equiv a \cdot \frac{\Delta U}{\Sigma U}$

beam


Size: 200x70 mm ${ }^{2}$


Shoe-box BPM:
Advantage: Very linear, low frequency dependence i.e. position sensitivity $\boldsymbol{S}$ is constant

Disadvantage: Large size, complex mechanics high capacitance

## Boundary Contribution $\Rightarrow$ FEM Calculation required

Boundary condition by the environment can significantly influence BPM properties
$\Rightarrow$ real properties have to be calculated numerically by Finite Element Method:
Examples are: CST-Studio (MAFIA), Comsol, HFFS

## General idea of FEM calculations:

$>$ Volume is divided in 3-dim meshes with typically $10^{6}$ to $10^{7}$ nodes
$>$ The beam is simulated by a traveling wave on a wire
$>$ Goal: Field distribution within the meshes
> The Maxwell equations are solved by iterative matrix inversion
$>$ Time domain: Propagation of source terms (here: Gaussian shaped pulse corresponding to 200 MHz bandwidth)
$>$ Frequency domain: e.g. eigenmodes

$>$ Output: time dependent signal, frequency dependences, S-parameters, field distribution etc.

## Optimization of Position Sensitivity

Simulation: Gaussian pulse travels on wire on different positions
$\rightarrow$ induced voltage calculated on matched output ports
$\rightarrow$ calculation of $\Delta \mathrm{U} / \Sigma \mathrm{U}$
Criteria of optimization: linearity, sensitivity, offset reduction, x-y plane independence


## Result:



Nearly perfect: $S_{x}=0.96 \% / \mathrm{mm}, \delta_{x}=-0.4 \mathrm{~mm}$ (ideal value $S_{x}=1.1 \% / \mathrm{mm}, \delta=0$ )

$$
\left.S_{y}=2.6 \% / \mathrm{mm}, \quad \delta_{y}=-0.04 \mathrm{~mm} \text { (ideal value } S_{y}=2.9 \% / \mathrm{mm}, \delta=0\right) \text { at } \mathbf{1} \mathbf{~ M H z}
$$

## Frequency Dependence of Position Sensitivity


$>$ Displacement sensitivity is nearly frequency independent only with separating guard rings
$>$ Sensitivity with separating rings is a factor of two larger as without ring.
> Capacitive cross talk spoils the sensitivity

## Plate-to-Plate Cross-Talk reduces Sensitivity

-Capacitive coubling determines position sensitivity
-Plate-to-plate cross talk caused by ceramic permittivity $\varepsilon=9.6$ resulting in high coupling capacitance between adjacent plates $\Rightarrow$ Insertion of the guard-ring between plates reduces cross talks by more than 10 dB



Remark: If guard rings are too large: E-field in-homogeneities occur

## Technical Realization of Shoe-Box BPM

Technical realization at HIT synchrotron of 46 m length for $7 \mathrm{MeV} / \mathrm{u} \rightarrow 440 \mathrm{MeV} / \mathrm{u}$ BPM clearance: 180x70 $\mathrm{mm}^{2}$, standard beam pipe diameter: 200 mm .


## Technical Realization of Shoe-Box BPM

Technical realization at HIT synchrotron of 46 m length for $7 \mathrm{MeV} / \mathrm{u} \rightarrow 440 \mathrm{MeV} / \mathrm{u}$ BPM clearance: $180 \times 70 \mathrm{~mm}^{2}$, standard beam pipe diameter: 200 mm .


## Other Types of diagonal cut BPM

Round type: cut cylinder
Same properties as shoe-box:


Other realization: Full metal plates
$\rightarrow$ No guard rings required
$\rightarrow$ but mechanical alignment more difficult

Wounded strips:
Same distance from beam and capacitance for all plates
But horizontal-vertical coupling.


# Beam Position Monitor: <br> Detector Principle, Hardware and Electronics 

Outline:
$>$ Signal generation $\rightarrow$ transfer impedance
> Consideration for capacitive shoe box BPM
> Consideration for capacitive button BPM
simple electro-static model, low $\beta$ effect, modification for synch. light source Comparison shoe box button BPM
$>$ Other BPM principles: stripline $\rightarrow$ traveling wave, inductive $\rightarrow$ wall current, cavity $\rightarrow$ resonator for dipole mode
$>$ Electronics for position evaluation
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## Button BPM for short Bunches

LINACs, e-synchrotrons: $100 \mathrm{MHz}<f_{r f}<3 \mathrm{GHz} \rightarrow$ bunch length $\approx$ BPM length $\rightarrow 50 \Omega$ signal path to prevent reflections
Button BPM with $50 \Omega \Rightarrow U_{i m}(t)=R \cdot \frac{A}{\beta c \cdot 2 \pi a} \cdot \frac{d I_{\text {beam }}}{d t}$
Example: LHC-type inside cryostat: $\varnothing 24 \mathrm{~mm}$, half aperture $a=25 \mathrm{~mm}, C=8 \mathrm{pF}$
$\Rightarrow f_{\text {cut }}=400 \mathrm{MHz}, Z_{t}=1.3 \Omega$ above $f_{\text {cut }}$


From C. Boccard (CERN)


EEin

## 2-dim Model for Button BPM

## 'Proximity effect': larger signal for closer plate

Ideal 2-dim case: Cylindrical pipe $\rightarrow$ image current density via 'image charge method' for 'pensile' beam:

$$
j_{i m}(\phi)=\frac{I_{\text {beam }}}{2 \pi a} \cdot\left(\frac{a^{2}-r^{2}}{a^{2}+r^{2}-2 a r \cdot \cos (\phi-\theta)}\right)
$$



Image current: Integration of finite BPM size: $I_{i m}=a \cdot \int_{-\alpha / 2}^{\alpha / 2} j_{i m}(\phi) d \phi$



## 2-dim Model for Button BPM

Ideal 2-dim model: Non-linear behavior and hor-vert coupling:
Sensitivity: $x=1 / S \cdot \Delta U / \Sigma U$ with $S[\% / \mathrm{mm}]$ or $[\mathrm{dB} / \mathrm{mm}]$ For this example: center part $S=7.4 \% / \mathrm{mm} \Leftrightarrow k=1 / S=14 \mathrm{~mm}$



Current density can also be calculated by Laplace equation for Fourier components
$I_{\text {beam }}=\left\langle I_{\text {beam }}\right\rangle+2\left\langle I_{\text {beam }}\right\rangle \cdot \sum_{n=1}^{\infty} A_{n} \cos \left(n \omega_{0} t\right) \quad$ for Gaussian bunches : $A_{n}=\exp \left(-n^{2} \omega^{2} \sigma_{t}^{2} / 2\right)$
In addition, frequency dependence can be calculated by this method.
P. Forck et al., GSI, February, 2009

## Position Measurement for Button BPM

Example LHC type: Measurement with movable $50 \Omega$ matched antenna:



From C. Boccard, C. Palau-Montava et al.(CERN).

## Estimation of finite Beam Size Effect for Button BPM

## Ideal 2-dim model:

Due to the non-linearity, the beam size enters in the position reading.

## Finite beam size:



Remark: For most LINACs: Linearity is less important, because beam has to be centered $\rightarrow$ correction as feed-forward for next macro-pulse.

## FEM Calculation for Button BPM simple Test Case

For realistic beam, 3-dim FEM calculations are required.

- Example: Button BPM at $\mathrm{r}=3 \mathrm{~cm}$ beam-pipe, flat, round $\varnothing 4 \mathrm{~cm}$
frequency $f_{r f}=150 \mathrm{MHz}$, effect for higher harmonics calculated



Nearly same result as ideal case!

## Low Velocity Effect: General Consideration

Simple Lorentz transformation of single point-like charge:
Lorentz boost and transformation of time: $\quad E_{\perp}(t)=\gamma E^{\prime}\left(t^{\prime}\right)$ and $t \rightarrow t^{\prime}$

$$
\text { E-field of a point-like charge: } \quad E_{\perp}(t)=\frac{e}{4 \pi \varepsilon_{0}} \cdot \frac{\gamma R}{\left[R^{2}+(\gamma \beta c t)^{2}\right]^{3 / 2}}
$$




## FEM Calculation of low $\beta$ Effect for p-LINAC



## Realization of Button BPM at LHC

Example LHC: $\varnothing 24 \mathrm{~mm}$, half aperture $a=25 \mathrm{~mm}$, installed inside cryostat Critically: $50 \Omega$ matching of button to standard feed-through.


From C. Boccard,
C. Palau-Montava et al.(CERN).

## Button BPM at Synchrotron Light Sources

The button BPM can be rotated by $45^{0}$
to avoid exposure by synchrotron light:
Frequently used at boosters for light sources

horizontal : $x=\frac{1}{S} \cdot \frac{\left(U_{1}+U_{4}\right)-\left(U_{2}+U_{3}\right)}{U_{1}+U_{2}+U_{3}+U_{4}}$
vertical: $\quad y=\frac{1}{S} \cdot \frac{\left(U_{1}+U_{2}\right)-\left(U_{3}+U_{4}\right)}{U_{1}+U_{2}+U_{3}+U_{4}}$

Example: Booster of ALS, Berkeley


## Button BPM at Synchrotron Light Sources

Due to synchrotron radiation, the button insulation might be destroyed
$\Rightarrow$ buttons only in vertical plane possible $\Rightarrow$ increased non-linearity


From S. Varnasseri, SESAME, DIPAC 2005

$>$ Beam position swept with 2 mm steps
$>$ Non-linear sensitivity and hor.-vert. coupling
$\Rightarrow$ At center $S_{x}=8.5 \% / \mathrm{mm}$ in this case
horizontal : $x=\frac{1}{S_{x}} \cdot \frac{\left(U_{2}+U_{4}\right)-\left(U_{1}+U_{3}\right)}{U_{1}+U_{2}+U_{3}+U_{4}}$
vertical: $y=\frac{1}{S_{y}} \cdot \frac{\left(U_{1}+U_{2}\right)-\left(U_{3}+U_{4}\right)}{U_{1}+U_{2}+U_{3}+U_{4}}$
Beam Position Monitors: Principle and Realization

## Button BPM at Synchrotron Light Sources


$>x \& y$ dependent polynomial fit possible


- Beam position swept with 2 mm steps
P. Forck et al., GSI, February, 2009

From S. Varnasseri, SESAME, DIPAC 2005
Beam Position Monitors: Principle and Realization

## Button BPM at Synchrotron Light Sources

2-dim electro-static simulation:


Result:
Distance $\mathrm{d}_{\mathrm{AB}}$ influences the sensitivity


Larger $\mathrm{d}_{\mathrm{AB}}$ has the effect:
$>$ higher sensitivity in x -direction
$>$ lower sensitivity in $y$-direction
$>$ linearity in influenced
$\Rightarrow$ Numerical optimization required

- Beam position swept with 2 mm steps $\qquad$ From S. Varnasseri, SESAME, DIPAC 2005
P. Forck et al., GSI, February, 2009


## Comparison Shoe-Box and Button BPM

|  | Shoe-Box BPM | Button BPM |
| :--- | :--- | :--- |
| Precaution | Bunches longer than BPM | Bunch length comparable to BPM |
| BPM length (typical) | 10 to 20 cm length per plane | $\varnothing 1$ to 5 cm per button |
| Shape | Rectangular or cut cylinder | Orthogonal or planar orientation |
| Bandwidth (typical) | 0.1 to 100 MHz | 100 MHz to 5 GHz |
| Coupling | $1 \mathrm{M} \Omega$ or $\approx 1 \mathrm{k} \Omega$ (transformer) | $50 \Omega$ |
| Cutoff frequency (typical) | $0.01 \ldots 10 \mathrm{MHz}(C=30 \ldots 100 \mathrm{pF})$ | $0.3 \ldots 1 \mathrm{GHz}(C=2 \ldots 10 \mathrm{pF})$ |
| Linearity | Very good, no x-y coupling | Non-linear, x-y coupling |
| Sensitivity | Good, care: plate cross talk | Good, care: signal matching |
| Usage | At proton synchrotrons, <br> $f_{r f}<10 \mathrm{MHz}$ | All electron acc., proton Linacs, <br> $f_{r f}>100 \mathrm{MHz}$ |



