Atomic Physics @ GSI (and FAIR)

All you (n)ever wanted to know about atomic physics with heavy ions

Harald Bräuning

SD / AP



Atomic Physics



AP @ GSI



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Highly Charged Heavy Ions

Hydrogen

simplest atomic system

best studied atomic system

energy difference between the 1s and 2s measured with **10**⁻¹⁴ precision by Laser spectroscopy

Heavy lons

extremely strong electric fields relativistic effects become significant QED effect increase with Z⁴ QED becomes more difficult to calculate simple few electron systems



Quantum Electro-Dynamics

'...my physics students don't understand it either. That is because I do not understand it. Nobody does.'

'The theory ... describes Nature as absurd from the point of view of common sense. ... So I hope you can accept Nature as She is – absurd.'

Richard P. Feynman: QED - The Strange Theory of Light and Matter



Quantum Electro-Dynamics

n=2 energy levels in hydrogen



QED contribution scales with Z⁴/n³

Higher Order QED Contributions



Measurements at the Electron Cooler



Measurements at the Electron Cooler

Gumberidze et al.: Phys. Rev. Lett. 94 (2005) 223001

From the $Ly \alpha_1$	From the K-RR	Mean value	Finite nuclear size	198.81
460.9 ± 2.5	454.9 ± 5.4	459.8 ± 2.3	Nuclear Recoil	0.46
The final result for	450 0 1 4 0		Nuclear Polarization	-0.19
the 1s Lamb shift	459.8 ± 4.2		VP (see Fig 2.1)	-88.60
520 510 490 490 490 400 450 440 430 420 1990 1992			SE (see Fig 2.1)	355.05
	scelerated Cooler	Decelerated lons: Cooler (our exp.)	SESE (see Fig 2.2)	-1.87
			VPVP (see Fig 2.2)	-0.97
			SEVP (see Fig 2.2)	1.14
			S(VP)E (see Fig 2.2)	0.13
			Total Lamb shift	463.95 ± 0.5
			Experiment	459.8 ± 4.2
	9941996199820002002YearExperiment not sensitive to hig order contributions		tive to higher	







Comparison: Ge(i) detector - crystal-spectrometer

Ge(i) pulse height $\epsilon = 10^{-4}$

crystal spectrometer $\epsilon = 10^{-8}$







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1. Beamtime with a 2d μ -strip detector: March 2006

H. Beyer, R. Reuschl, et al.

raw x-ray image

x-ray image in coincidence with down-charged ion (electron capture)

B x-ray image (10 keV to 100 keV) x-ray image 0 keV to 100 keV) + time condition x-ray image (55 keV to 65 keV) M + time condition × [mm] ray image to 100 ke x-ray image (10 keV to 100 keV) + time condition x-ray image (55 keV to 65 keV) + time condition E × (mm) W x-ray image keV to 100 keV) (mm) x-ray image 10 keV to 100 keV) + time condition E x-ray image (55 keV to 65 keV) + time condition m 10 × (mm)

x-ray in image in coincidence with down-charged ion and with preselected x-ray energies (58-65 keV)



1. Beamtime with a 2d $\mu\text{-strip}$ detector: March 2006

H. Beyer, R. Reuschl, et al.





Problems:

resolution determined by width of detector strips detectors with smaller strips not realistic

Current Development:

sub-pixel resolution via pulse shape analysis



E = hv

Photoabsorption

Periodic Potential







photon energy \rightarrow excitation energy

 $E = hv = h\frac{v}{d}$

(non relativistic)

kinetic energy \rightarrow excitation energy



planar channeling

$$E_{trans} = \frac{h v \gamma}{a} \left(\sqrt{2} k \cos \theta + l \sin \theta \right)$$

a: lattice constant; k,l,m: Miller-Indices

Tunable source of virtual photons:

coarse tuning: beam energy

fine tuning: crystal orientation



Condition for 1s - 2p Transitions

Resonance Energy of 1s-2p in Si[110]





5 axis goniometer



 $0.4 \ \mu rad$ angular resolution

Made in Japan

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Made in Germany

GSI



Beam Diagnostics



detection of the $2p \rightarrow 2s$ transition

Observation angles: $\pm 43.4 \text{ deg and } \pm 32.8 \text{ deg}$

 \mathbf{E}_{lab} = 6.97 KeV and 6.25 keV for the 4.4 6 keV Uranium transition



Detected x-rays as function of crystal orientation

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IPN Lyon D. Dauvergne

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Tomsk University K. Klimova Y. Pivovarov





Radiative Electron Capture



Photoionization



H. Stobbe, Ann. Phys. 7 (1931) 661

Radiative Electron Capture

REC

Photoionization



GSI

Radiative Electron Capture



Relativistic effects decrease the linear polarization

For high energies a "cross-over" effect can be observed Ann. Phys. 9 (1931) 21

A.Surzhykov et al., Phys. Lett. A 289 (2001) 213; J. Eichler et al., Phys. Rev. A 65 (2002) 052716

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ESR Internal Target Section



Compton - Polarimetry

Compton - Scattering

Linearly Polarized Radiation



Compton - Polarimetry



Micro – Strip Detectors

New micro-strip semiconductor detectors

Si(Li) or Ge(i) energy resolution timing 2D (3D) position sensitivity multi-hit capability

single crystal for Compton scattering and absorption



Micro – Strip Detectors

2D position sensitive Si(Li) detector

64 x 64 mm² active area: crystal thickness: 7 mm number of strips: 32 + 32 pitch: 2mm



active area: $64 \times 64 \text{ mm}^2$ crystal thickness: 11 mm number of strips: 48 + 128 1167µm and 250µm pitch:



D+-

contact



GSĬ

germanium detector



64

mm

Compton - Polarimetry

Images of Compton scattering distributions for well defined scattering angles

 $\mathsf{Energy} \leftrightarrow \mathsf{Angle}$

$$E' = \frac{E}{1 + \frac{E}{m_e c^2} (1 - \cos \theta_c)}$$

Test of the Polarization Sensitivity at the ESRF Synchrotron Facility using 100% Linearly Polarized Radiation



Results of a test beamtime at the ESRF with 98% linearly polarized x-rays

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Experimental Set-Up



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96.6 MeV/u U⁹²⁺ + H₂

X-ray spectrum after electron capture



96.6 MeV/u U⁹²⁺ + H₂





96.6 MeV/u U⁹²⁺ + H₂

K-REC



Doktorarbeit: S. Hess

Future Applications in Beam Diagnostics

REC as a 'probe' for measuring the ion spin-polarization Surzhykov et al., PRL 94 (2005) 203202

spin-polarized, heavy ions (Z > 54)

parity non conservation studies permanent electric dipole momont spin effects in collisions





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Future Applications in Beam Diagnostics

REC as a 'probe' for measuring the ion spin-polarization Surzhykov et al.,



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Future Applications in Beam Diagnostics

REC as a 'probe' for measuring the ion spin-polarization Surzhykov et al., PRL 94 (2005) 203202

Stokes parameter P_1 is polarization independent

Stokes parameter P₂ is strongly dependent on degree of polarization

spin polarization leads to a rotation of the polarization plane



unpolarized ion beam: $P_2 = 0$

polarized ion beam: $P_2 \neq 0$

$$\tan 2\Psi = \frac{P_2}{P_1}$$

A. Surzhykov et al., PRL 94 (2005) 203202

The Crew

rov, R. Reuschl, D. Protic, U. Spillmann, Th. Stöhlker, M. Trassinelli, S. Trotsenko, G. Webe

Theory

J. Eichler, S. Fritzsche, A. Ichihara, A. Surzhykov

Theoretische Physik, HMI-Berlin, Germany JAERI, Japan University of Heidelberg, Germany

... and many more







